

J.D. Jackson Problem 5.27

Josh Orndorff
admin@joshorndorff.com

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We'll begin by finding the magnetic field in the three regions. This calculation is somewhat simple because the problem's symmetry allows the use of Ampere's law.

$$\mathbf{B}_{\text{in}} = \frac{\mu I r}{2\pi b^2} \hat{\phi} \quad (1)$$

$$\mathbf{B}_{\text{mid}} = \frac{\mu_0 I}{2\pi r} \hat{\phi} \quad (2)$$

$$\mathbf{B}_{\text{out}} = \mathbf{0} \quad (3)$$

By equation (5.148) and a simplified version of (5.152) we see that

$$L = \frac{1}{I^2} \int \frac{B^2}{\mu} dV \quad (4)$$

$$L = \frac{1}{I^2} \int_0^{2\pi} \int_0^l \left[\int_0^b \frac{\mu_0^2 I^2 r^2}{4\pi^2 \mu b^4} r dr + \int_b^a \frac{\mu_0^2 I^2}{4\pi^2 \mu_0 r^2} r dr + 0 \right] dz d\theta \quad (5)$$

$$\boxed{\frac{L}{l} = \frac{\mu_0^2}{2\pi} \left[\frac{1}{4\mu} + \frac{1}{\mu_0} \ln \frac{a}{b} \right]} \quad (6)$$

In the case that the inner conductor is hollow, the B -field in the inner-most region becomes zero, and the Inductance per unit length simplifies

$$\boxed{\frac{L}{l} = \frac{\mu_0}{2\pi} \ln \frac{a}{b}} \quad (7)$$

If we return to the more common definition of inductance (**not** per unit length) we get

$$\xi = -\frac{d}{dt} \iint \mathbf{B} \cdot d\mathbf{A} \quad (8)$$

$$= -l \frac{d}{dt} \int_b^a \frac{\mu I}{2\pi r} dr \quad (9)$$

$$= -\frac{l\mu_0}{2\pi} \frac{dI}{dt} \ln \frac{a}{b} \quad (10)$$

$$\boxed{L = \frac{-\xi}{\frac{dI}{dt}} = \frac{l\mu_0}{2\pi} \ln \frac{a}{b}} \quad (11)$$